# Asymptotics of the Stirling numbers of the first kind revisited: A saddle point approach 

Guy Louchard

Université Libre de Bruxelles, Dép. d'Informatique, CP 212, B-1050 Bruxelles, Belgium. louchard@ulb.ac.be
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Using the saddle point method, we obtain from the generating function of the Stirling numbers of the first kind $\left[\begin{array}{c}n \\ j\end{array}\right]$ and Cauchy's integral formula, asymptotic results in central and non-central regions. In the central region, we revisit the celebrated Goncharov theorem with more precision. In the region $j=n-n^{\alpha}, \alpha>1 / 2$, we analyze the dependence of $\left[\begin{array}{c}n \\ j\end{array}\right]$ on $\alpha$.

Keywords: Stirling numbers, saddle point method.

To Philippe

## 1 Introduction

Let $\left[\begin{array}{l}n \\ j\end{array}\right]$ be the Stirling number of the first kind (unsigned version). Their generating function is given by

$$
\phi_{n}(z)=\prod_{0}^{n-1}(z+i)=\frac{\Gamma(z+n)}{\Gamma(z)}, \quad \phi_{n}(1)=n!
$$

In the sequel all asymptotics are meant for $n \rightarrow \infty$.
An asymptotic expansion for $j=\mathcal{O}(1)$ is given in Wilf [14], which has been extended to the range $j=\mathcal{O}(\ln n)$ by Hwang [6]. The generalized Stirling numbers have been considered by Tsylova [13] and Chelluri et al. [2]. The $q$-Stirling numbers are studied in Kyriakoussis and Vamvakari [9].

In Sec 2, we revisit the asymptotic expansions in the central region and in $\operatorname{Sec} 3$, we analyse the noncentral region $j=n-n^{\alpha}, \quad \alpha>1 / 2$. We use Cauchy's integral formula and the saddle point method.

## 2 Central region

Consider the random variable $J_{n}$, with probability distribution

$$
\begin{aligned}
& \mathbb{P}\left[J_{n}=j\right]=Z_{n}(j), \\
Z_{n}(j):= & \frac{\left[\begin{array}{l}
n \\
j
\end{array}\right]}{n!} .
\end{aligned}
$$

The mean and variance are given by

$$
\begin{aligned}
M & :=\mathbb{E}\left(J_{n}\right)=\sum_{0}^{n-1} \frac{1}{1+i}=H_{n}=\psi(n+1)+\gamma \\
\sigma^{2} & :=\mathbb{V}\left(J_{n}\right)=\sum_{0}^{n-1} \frac{i}{(1+i)^{2}}=\psi(1, n+1)+\psi(n+1)-\frac{\pi^{2}}{6}+\gamma
\end{aligned}
$$

where $\psi(x)$ is the digamma function, $\psi(k, x)$ is the $k$ th polygamma function, and

$$
\begin{aligned}
M & \sim \ln (n)+\gamma+\frac{1}{2 n}+\mathcal{O}\left(\frac{1}{n^{2}}\right) \\
\sigma^{2} & \sim \ln (n)-\frac{\pi^{2}}{6}+\gamma+\frac{3}{2 n}+\mathcal{O}\left(\frac{1}{n^{2}}\right)
\end{aligned}
$$

It is convenient to set

$$
A_{n}:=\ln (n)-\frac{\pi^{2}}{6}+\gamma=\ln \left(n e^{\gamma-\pi^{2} / 6}\right)
$$

and to consider all our next asymptotics $(n \rightarrow \infty)$ as functions of $A_{n}$. Of course, all asymptotics can be reformulated in terms of $\ln (n)$.

We have

$$
\begin{aligned}
M & \sim A_{n}+\frac{\pi^{2}}{6}+\mathcal{O}\left(\frac{1}{n}\right) \\
\sigma^{2} & \sim A_{n}+\mathcal{O}\left(\frac{1}{n}\right)
\end{aligned}
$$

A celebretated central limit theorem of Goncharov says that

$$
J_{n} \sim \mathcal{N}(M, \sigma)
$$

where $\mathcal{N}$ is the Gaussian distribution, with a rate of convergence $\mathcal{O}(1 / \sqrt{\ln (n)})$. This can also be deduced from the Quasi-Power theorem of Hwang [7], [8].

In this Section, we want to obtain a more precise local limit theorem for $J_{n}$ in terms of $x:=\frac{J_{n}-M}{\sigma}$ and $A_{n}$. Actually, we obtain the following theorem, where we use $B_{n}:=\sqrt{A_{n}}$ to simplify the expressions.

## Theorem 2.1

$$
\begin{aligned}
Z_{n}(j) & \sim \frac{1}{\sqrt{2 \pi} B_{n}} e^{-x^{2} / 2} . \\
& \cdot\left[1+\frac{x^{3} / 6-x / 2}{B_{n}}+\frac{3 x^{2} / 8-x^{4} / 6-1 / 12+x^{6} / 72}{B_{n}^{2}}\right. \\
& \left.+\frac{-\pi^{2} x^{3} / 18+37 x^{5} / 240-355 x^{3} / 144+x / 8-x^{7} / 48+x^{9} / 1296+\pi^{2} x / 6-\zeta(3) x+\zeta(3) x^{3} / 3}{B_{n}^{3}}+\ldots\right] .
\end{aligned}
$$

Proof: By Cauchy's theorem,

$$
\begin{aligned}
Z_{n}(j) & =\frac{1}{2 \pi \mathbf{i}} \int_{\Omega} \frac{\phi_{n}(z)}{z^{j+1} n!} d z \\
& =\frac{1}{2 \pi \mathbf{i}} \int_{\Omega} e^{S(z)} d z
\end{aligned}
$$

where $\Omega$ is inside the analyticity domain of the integrand and encircles the origin and

$$
\left.S(z)=S_{1}(z)+S_{2}(z), \quad S_{1}(z)=\sum_{i=0}^{n-1} \ln (z+i)\right)-\ln (n!), \quad S_{2}(z)=-(j+1) \ln (z)
$$

We will use the Saddle point method (for a good introduction to this method, see Flajolet and Sedgewick [3], ch.VIII). Set

$$
S^{(i)}:=\frac{d^{i} S}{d z^{i}}
$$

These derivatives can be expressed in terms of $\psi(k, z+n)$ and $\psi(k, z)$.
First we must find the solution of

$$
\begin{equation*}
S^{(1)}(\tilde{z})=0 \tag{1}
\end{equation*}
$$

with smallest module.
Set $\tilde{z}:=z^{*}-\varepsilon$, where $z^{*}=\lim _{n \rightarrow \infty} \tilde{z}$. Here, it is easy to check that $z^{*}=1$. Set $j=M+x \sigma, x$ fixed and $B_{n}:=\sqrt{A_{n}}$.

This leads, to first order (keeping only the $\varepsilon$ term in (1)), to

$$
\varepsilon:=\frac{-x}{B_{n}}+\frac{x^{2}-1}{B_{n}^{2}}+\ldots+\frac{1}{n}\left(\frac{3 x}{4 B_{n}^{3}}+\ldots\right)+\mathcal{O}\left(\frac{1}{n^{2} B_{n}^{4}}\right) .
$$

This shows that, asymptotically, $\varepsilon$ is given by a series of powers of $n^{-1}$, where each coefficient is given by a series of powers of $B_{n}^{-1}$. To obtain more precision, we set again $j=M+x \sigma$, expand in powers of $n^{-1}$, and equate each coefficient to 0 . . This leads to (here and in the following, we provide only a few terms but Maple knows more)

$$
\varepsilon=\frac{-x}{B_{n}}-\frac{1}{B_{n}^{2}}+\frac{0}{B_{n}^{3}}+\ldots+\frac{1}{n}\left(\frac{3 x}{4 B_{n}^{3}}+\frac{x^{2}+3 / 2}{B_{n}^{4}}+\ldots\right)+\mathcal{O}\left(\frac{1}{n^{2} B_{n}^{4}}\right)
$$

We have, with $\tilde{z}:=z^{*}-\varepsilon=1-\varepsilon$,

$$
Z_{n}(j)=\frac{1}{2 \pi \mathbf{i}} \int_{\Omega} \exp \left[S(\tilde{z})+S^{(2)}(\tilde{z})(z-\tilde{z})^{2} / 2!+\sum_{l=3}^{\infty} S^{(l)}(\tilde{z})(z-\tilde{z})^{l} / l!\right] d z
$$

Note that the linear term vanishes. Set $z=\tilde{z}+\mathbf{i} \tau$. This gives

$$
\begin{equation*}
Z_{n}(j)=\frac{1}{2 \pi} \exp [S(\tilde{z})] \int_{-\infty}^{\infty} \exp \left[S^{(2)}(\tilde{z})(\mathbf{i} \tau)^{2} / 2!+\sum_{l=3}^{\infty} S^{(l)}(\tilde{z})(\mathbf{i} \tau)^{l} / l!\right] d \tau \tag{2}
\end{equation*}
$$

The justification of the integration procedure is given in the Appendix. Let us first analyze $S(\tilde{z})$. We obtain

$$
\begin{aligned}
S(\tilde{z}) & =-x^{2} / 2+\frac{x^{3} / 6-x}{B_{n}}+\frac{-x^{4} / 12+x^{2} / 2-1 / 2}{B_{n}^{2}} \\
& +\frac{-x^{3} / 3+x^{5} / 20+x / 2-\pi^{2} x^{3} / 18+\zeta(3) x^{3} / 3}{B_{n}^{3}}+\ldots+\mathcal{O}\left(\frac{1}{n B_{n}^{3}}\right) .
\end{aligned}
$$

Also,

$$
\begin{aligned}
S^{(2)}(\tilde{z}) & =B_{n}^{2}-B_{n} x-1+x^{2}+\ldots, \\
S^{(3)}(\tilde{z}) & =-2 B_{n}^{2}+4 B_{n} x-\pi^{2} / 3+2 \zeta(3)-6 x^{2}+4+\ldots, \\
S^{(4)}(\tilde{z}) & =6 B_{n}^{2}-18 B_{n} x+36 x^{2}-18+\pi^{2}-\pi^{4} / 15+\ldots, \\
S^{(l)}(\tilde{z}) & =\mathcal{O}\left(B_{n}^{2}\right), l \geq 5 .
\end{aligned}
$$

We need these many terms in the following. Note that, with $z=\tilde{z} e^{\mathbf{i} \theta}$, this leads to

$$
\begin{equation*}
S^{(2)}(\tilde{z}) \frac{(z-\tilde{z})^{2}}{2} \sim-\frac{1}{2} \ln (n) \theta^{2} \tag{3}
\end{equation*}
$$

We can now compute (2), for instance by using the classical trick of setting

$$
S^{(2)}(\tilde{z})(i \tau)^{2} / 2!+\sum_{l=3}^{\infty} S^{(l)}(\tilde{z})(i \tau)^{l} / l!=-u^{2} / 2
$$

Computing $\tau$ as a truncated series in $u$, this gives, by inversion,

$$
\tau=\left[u\left(1+x /\left(2 B_{n}\right)+\ldots\right)+u^{2}\left(\mathbf{i} /\left(3 B_{n}\right)+\ldots\right)+u^{3}\left(-1 /\left(36 B_{n}^{2}\right)+\ldots\right)\right] / B_{n}+\ldots
$$

Setting $d \tau=\frac{d \tau}{d u} d u$, expanding w.r.t. $B_{n}$ and integrating on $[u=-\infty . . \infty]$, this gives

$$
\frac{1}{\sqrt{2 \pi} B_{n}}\left[1+\frac{x}{2 B_{n}}+\frac{5 / 12-x^{2} / 8}{B_{n}^{2}}+\frac{x\left(8 \pi^{2}-10-93 x^{2}-48 \zeta(3)\right)}{48 B_{n}^{3}}+\ldots\right]
$$



Fig. 1: Comparison between $Z_{n}(j)$ and $\frac{1}{\sqrt{2 \pi} \sigma} \exp \left[-\left(\frac{j-M}{\sigma}\right)^{2} / 2\right], n=3000$
Finally (2) leads to

$$
\begin{aligned}
Z_{n}(j) & \sim \frac{1}{\sqrt{2 \pi} B_{n}} e^{-x^{2} / 2} . \\
& \cdot \exp \left[\frac{x^{3} / 6-x}{B_{n}}+\frac{-x^{4} / 12+x^{2} / 2-1 / 2}{B_{n}^{2}}+\frac{-x^{3} / 3+x^{5} / 20+x / 2-\pi^{2} x^{3} / 18+\zeta(3) x^{3} / 3}{B_{n}^{3}}+\ldots\right] . \\
& \cdot\left[1+\frac{x}{2 B_{n}}+\frac{5 / 12-x^{2} / 8}{B_{n}^{2}}+\frac{x\left(8 \pi^{2}-10-93 x^{2}-48 \zeta(3)\right)}{48 B_{n}^{3}}+\ldots\right]
\end{aligned}
$$

or

$$
\begin{aligned}
Z_{n}(j) & \sim R_{1}, \\
R_{1} & =\frac{1}{\sqrt{2 \pi} B_{n}} e^{-x^{2} / 2} . \\
& \cdot\left[1+\frac{x^{3} / 6-x / 2}{B_{n}}+\frac{3 x^{2} / 8-x^{4} / 6-1 / 12+x^{6} / 72}{B_{n}^{2}}\right. \\
& \left.+\frac{-\pi^{2} x^{3} / 18+37 x^{5} / 240-355 x^{3} / 144+x / 8-x^{7} / 48+x^{9} / 1296+\pi^{2} x / 6-\zeta(3) x+\zeta(3) x^{3} / 3}{B_{n}^{3}}+\ldots\right] .
\end{aligned}
$$

For $n=3000$, a comparison between $Z_{n}(j)$ and $\frac{1}{\sqrt{2 \pi} \sigma} \exp \left[-\left(\frac{j-M}{\sigma}\right)^{2} / 2\right]$ is given in Figure 1


Fig. 2: $Z_{n}(j) /\left[\frac{1}{\sqrt{2 \pi} \sigma} \exp \left[-\left(\frac{j-M}{\sigma}\right)^{2} / 2\right]\right]$, color=red, $Z_{n}(j) / R_{1}$, color=blue, $n=3000$

Of course, only few values of $j$ are significant and also the quality of the Gaussian is low, all asymptotic expressions depend actually on powers of $A_{n}^{-1}$, but $A_{n}$ is not large.
A comparison of $Z_{n}(j) /\left[\frac{1}{\sqrt{2 \pi} \sigma} \exp \left[-\left(\frac{j-M}{\sigma}\right)^{2} / 2\right]\right]$ with $Z_{n}(j) / R_{1}$, with 2 terms in $R_{1}$, is given in Figure 2

The precision of $R_{1}$ is of order $10^{-2}$. Using 3 terms in $R_{1}$ leads to a less good result: $A_{n}$ is not large enough to take advantage of the $A_{n}^{-3 / 2}$ term: $A_{n}=6.94$ here, we deal with asymptotic series, not necessarily convergent ones. More terms can be computed in $R_{1}$ (which is almost automatic with Maple).

## 3 Large deviation, $j=n-n^{\alpha}, \quad 1>\alpha>1 / 2$

The case $j=\mathcal{O}(n)$ is analyzed in Timashev [12]. But he obtains a series of powers of $n^{-1}$, determined by a power series of a certain function that depends on the solution of a given non-linear differential equation of the first order. The coefficients obey some linear recurrence relations in the complex plane. The case $j=n-c, c$ constant, is considered in Grünberg [5]. As previous work for the case $j=n-n^{\alpha}$, let us mention Bender [1], Temme [11], Moser and Wyman [10] (see also the comments by Odlyzko in [4], p.1182). They all use, explicitly or not, the Saddle point method. For $\alpha<1 / 2$, Moser and Wyman (6.9) give an explicit asymptotic expression. For $\alpha>1 / 2$, they first compute in (4.52) the numerical solution $z n$ of $S^{\prime}(z n)=0$ and give in (4.51) an asymptotic expression. This is rather precise: for $n=50$, this gives a precision of order $10^{-4}$. [1] and [11] also compute numerically $z n$.

However, all these results do not shed light on the dependence of $\left[z^{j}\right] \phi(z)$ on $n^{\alpha}$. This is what we want
to explicit in this Section. It appears that the range $\alpha>1 / 2$ is more delicate than the other range.
Recall that

$$
\phi_{n}(z)=\prod_{0}^{n-1}(z+i)=\frac{\Gamma(z+n)}{\Gamma(z)}
$$

We have

$$
G_{n}(z):=\frac{\Gamma(z+n)}{\Gamma(z) z^{j+1}}=\exp [S(z)]
$$

with

$$
S(z)=S_{1}(z)+S_{2}(z), S_{1}(z)=\sum_{0}^{n-1} \ln (z+i), S_{2}(z)=-(j+1) \ln (z) .
$$

We first compute $\tilde{z}$ such that

$$
\begin{equation*}
S^{\prime}(\tilde{z})=0 \tag{4}
\end{equation*}
$$

We have

$$
S^{\prime}(z)=\psi(z+n)-\psi(z)-\frac{j+1}{z}
$$

Similarly (we need these expressions later on)

$$
\begin{aligned}
& S^{(2)}(z)=\psi(1, z+n)-\psi(1, z)+\frac{j+1}{z^{2}} \\
& S^{(k)}(z)=\psi(k-1, z+n)-\psi(k-1, z)+(-1)^{k}(k-1)!\frac{j+1}{z^{k}}
\end{aligned}
$$

Some experiments with some values for $\alpha$ ( $\alpha=5 / 8$ is a good choice) show that $\tilde{z}$ must be a combination of $x=n^{\alpha}$ and $y=n^{1-\alpha}$ and $x \gg y \gg 1$. Note that both $x$ and $y$ are large. We will derive series of powers of $x^{-1}$, where each coefficient is a series of powers of $y^{-1}$. We obtain the following theorem
Theorem 3.1

$$
\begin{aligned}
{\left[z^{j}\right] \phi_{n}(z) \sim } & \frac{1}{\sqrt{2 \pi}} \frac{y^{2} \sqrt{x}}{2} \exp \left[x \left[1-\ln (2)+2 \ln (y)+\ln (x)-\frac{2}{3 y}-\frac{2}{9 y^{2}}-\frac{44}{405 y^{3}}-\frac{26}{405 y^{4}}+\frac{40}{27 y^{5}}\right.\right. \\
& \left.\left.+\frac{179968}{18225 y^{6}}+\frac{4727552}{127575 y^{7}}+\frac{3436796}{32805 y^{8}}+\frac{5492621728}{22143375 y^{9}}+\ldots\right]+\ln (2)-2 \ln (y)-\ln (x)\right] \\
& \cdot\left[1-\frac{3}{3 y}-\frac{1}{18 y^{2}}-\frac{1}{30 y^{3}}+\frac{17207}{3240 y^{4}}+\ldots+\frac{1}{x}\left(-\frac{1}{12}+\frac{1}{36 y}-\frac{35}{216 y^{2}}+\frac{15029}{3240 y^{3}}+\ldots\right)\right. \\
& \left.+\frac{1}{x^{2}}\left(\frac{1}{288}-\frac{1}{864 y}+\frac{3527}{5184 y^{2}}+\ldots\right)+\mathcal{O}\left(\frac{1}{x^{3}}\right)\right]
\end{aligned}
$$

Proof: Let us summarize the different steps of the proof. First we compute $\tilde{z}$ and $S(\tilde{z})$ as $S(\tilde{z})=T_{1} T_{2}$, where $T_{1}$ is the dominant term and $T_{2}$ is a series of powers of $x^{-1}$, where each coefficient is a series of powers of $y^{-1}$. We expand $T_{3}:=e^{T_{2}}$. Next the integration procedure leads to $\frac{y^{2} \sqrt{x}}{2} T_{4}$, where $T_{4}$ is again a series of powers of $x^{-1}$, where each coefficient is a series of powers of $y^{-1}$. We set $T_{5}:=\frac{1}{\sqrt{2 \pi}} \frac{y^{2} \sqrt{x}}{2} e^{T_{1}}$. Finally, we obtain

$$
\begin{equation*}
\left[z^{j}\right] \phi_{n}(z) \sim T_{5} T_{3} T_{4} \tag{5}
\end{equation*}
$$

The first terms in the asymptotics of $\tilde{z}$ are easy to compute: set $\tilde{z}=n \beta$. Equation (4) leads to

$$
\psi(n(1+\beta))-\psi(n \beta)=\frac{1}{\beta}-\frac{1}{y \beta}+\frac{1}{n \beta} .
$$

But $\psi(n) \sim \ln (n)$. So we have

$$
\ln \left(1+\frac{1}{\beta}\right) \sim \frac{1}{\beta}-\frac{1}{y \beta}+\frac{1}{n \beta}
$$

or

$$
\frac{1}{\beta}-\frac{1}{2 \beta^{2}} \sim \frac{1}{\beta}-\frac{1}{y \beta}
$$

or $\beta \sim \frac{y}{2}$.
More generally, we have
$\beta=\frac{y}{2}\left[1+\frac{a_{1}}{y}+\frac{a_{2}}{y^{2}}+\frac{a_{3}}{y^{3}}+\ldots+\frac{1}{x}\left(1+\frac{b_{1}}{y}+\frac{b_{2}}{y^{2}}+\ldots\right)+\frac{1}{x^{2}}\left(1+\frac{c_{1}}{y}+\frac{c_{2}}{y^{2}}+\ldots\right)+\mathcal{O}\left(\frac{1}{x^{3}}\right)\right]$.
Note that $\frac{1}{y^{3}}$ can be of the same order than $\frac{1}{x}$, see below.
By bootstrapping, we obtain (we give the first terms)

$$
\begin{align*}
\tilde{z} & =\frac{n y}{2}\left[1-\frac{4}{3 y}+\frac{2}{9 y^{2}}+\frac{8}{135 y^{3}}+\frac{8}{405 y^{4}}+\frac{16}{1701 y^{5}}+\frac{232}{45525 y^{6}}+\frac{64}{18225 y^{7}}+\ldots\right. \\
& +\frac{1}{x}\left[1-\frac{1}{y}+\frac{4}{9 y^{2}}-\frac{16}{135 y^{3}}+\ldots\right] \\
& +\frac{1}{x^{2}}\left[1-\frac{1}{y}+\frac{0}{y^{2}}+\ldots\right] \\
& +\frac{1}{x^{3}}[1+\ldots] \\
& \left.+\mathcal{O}\left(\frac{1}{x^{4}}\right)\right] \tag{6}
\end{align*}
$$

Note that the choice of dominant terms in the bracket of (6) depends on $\alpha$. For instance, for $\alpha=3 / 4$, the dominant terms (in decreasing order) are

$$
1, \frac{1}{y}, \frac{1}{y^{2}},\left\{\frac{1}{x}, \frac{1}{y^{3}}\right\},\left\{\frac{1}{x y}, \frac{1}{y^{4}}\right\},\left\{\frac{1}{x y^{2}}, \frac{1}{y^{5}}\right\},\left\{\frac{1}{x^{2}}, \frac{1}{x y^{3}}, \frac{1}{y^{6}}\right\}, \ldots
$$

Now we must compute $S(\tilde{z})$ and its asymptotics. First we compute $\ln (\tilde{z}+i)$, take the asymptotics wrt $x$, sum on $i$, and again take the asymptotics wrt $x$ (recall that $n=x y$ ). this leads to

$$
\begin{aligned}
S_{1}(\tilde{z}) & =x\left[(-\ln (2)+2 \ln (y)+\ln (x)) y-\frac{1}{3}+\frac{4}{405 y^{2}}+\frac{2}{405 y^{3}}+\ldots\right]+y-\frac{2}{3}-\frac{2}{3 y}-\frac{49}{135 y^{2}}+\ldots \\
& +\frac{1}{x}\left(\frac{y}{2}+\frac{1}{6 y}+\ldots\right)+\frac{1}{x^{2}}\left(\frac{y}{3}+\ldots\right)+\mathcal{O}\left(\frac{y}{x^{3}}\right) .
\end{aligned}
$$



Fig. 3: $z n / \tilde{z}, n=500$, as function of $j$, full range


Fig. 4: $z n / \tilde{z}, n=500$, as function of $j$, restricted range

Here we provide only a few terms but Maple knows more. Next

$$
\begin{aligned}
S_{2}(\tilde{z}) & =x\left[(\ln (2)-2 \ln (y)-\ln (x)) y+\frac{4}{3}-\ln (2)+2 \ln (y)+\ln (x)-\frac{2}{3 y}-\frac{94}{405 y^{2}}+\ldots\right] \\
& -y+\frac{2}{3}+\ln (2)-2 \ln (y)-\ln (x)+\frac{1}{y}+\frac{94}{135 y^{2}}+\ldots \\
& +\frac{1}{x}\left(\frac{y}{2}+\frac{1}{6 y}+\ldots\right) \\
& +\frac{1}{x^{2}}\left(\frac{y}{3}+\ldots\right) \\
& +\mathcal{O}\left(\frac{y}{x^{3}}\right)
\end{aligned}
$$

So, finally

$$
\begin{aligned}
S(\tilde{z}) & =S_{1}(\tilde{z})+S_{2}(\tilde{z}) \sim x\left[1-\ln (2)+2 \ln (y)+\ln (x)-\frac{2}{3 y}+\ldots\right] \\
& +\ln (2)-2 \ln (y)-\ln (x)+\frac{1}{3 y}+\frac{1}{3 y^{2}}+\ldots \\
& +\frac{1}{x}\left(-\frac{1}{2}+\frac{1}{3 y}-\frac{1}{2 y^{2}}+\ldots\right) \\
& +\frac{1}{x^{2}}\left(-\frac{1}{6}+\frac{19}{18 y^{2}} \ldots\right) \\
& +\mathcal{O}\left(\frac{1}{x^{3}}\right)
\end{aligned}
$$

Now we split $S(\tilde{z})$ into two parts:

$$
\begin{aligned}
T_{1} & =x\left[1-\ln (2)+2 \ln (y)+\ln (x)-\frac{2}{3 y}+\ldots\right]+\ln (2)-2 \ln (y)-\ln (x) \\
T_{2} & =\frac{1}{3 y}+\frac{1}{3 y^{2}}+\ldots \\
& +\frac{1}{x}\left(-\frac{1}{2}+\frac{1}{3 y}-\frac{1}{2 y^{2}}+\ldots\right) \\
& +\frac{1}{x^{2}}\left(-\frac{1}{6}-\frac{17}{18 y^{2}} \ldots\right) \\
& +\mathcal{O}\left(\frac{1}{x^{3}}\right)
\end{aligned}
$$

Note that the dominant term of $T_{1}$ is given by

$$
\begin{equation*}
T_{1} \sim(2-\alpha) n^{\alpha} \ln (n) \tag{7}
\end{equation*}
$$

We obtain

$$
\exp (S(\tilde{z}))=e^{T_{1}} e^{T_{2}}=e^{T_{1}} T_{3}
$$

with

$$
\begin{aligned}
T_{3} & =e^{T_{2}}=1+\frac{1}{3 y}+\frac{7}{18 y^{2}}+\frac{89}{270 y^{3}}+\frac{18263}{3240 y^{4}}+\frac{98009}{3240 y^{5}}+\frac{9517337}{97200 y^{6}}+\frac{491504273}{2041200 y^{7}}+\ldots \\
& +\frac{1}{x}\left(-\frac{1}{2}+\frac{1}{6 y}-\frac{7}{12 y^{2}}+\frac{2311}{540 y^{3}}+\frac{112469}{6480 y^{4}}+\frac{5137}{144 y^{5}}+\ldots\right) \\
& +\frac{1}{x^{2}}\left(-\frac{1}{24}-\frac{13}{72 y}-\frac{557}{932 y^{2}}+\ldots\right) \\
& +\mathcal{O}\left(\frac{1}{x^{3}}\right)
\end{aligned}
$$

Here we have given all terms compatible with the expansion (6). Also, with more precision,

$$
\begin{aligned}
T_{1} & =x\left[1-\ln (2)+2 \ln (y)+\ln (x)-\frac{2}{3 y}-\frac{2}{9 y^{2}}-\frac{44}{405 y^{3}}-\frac{26}{405 y^{4}}+\frac{40}{27 y^{5}}\right. \\
& \left.+\frac{179968}{18225 y^{6}}+\frac{4727552}{127575 y^{7}}+\frac{3436796}{32805 y^{8}}+\frac{5492621728}{22143375 y^{9}}+\ldots\right] \\
& +\ln (2)-2 \ln (y)-\ln (x) .
\end{aligned}
$$

Now we must consider $S^{(k)}(\tilde{z})$. By direct expansion, this gives the following expressions (again we provide only the first few terms). We must use up to six derivatives to get a sufficient precision (of order $x^{-2}$ ) in the Saddle integrals.

$$
\begin{align*}
S^{(2)}(\tilde{z}) & =\frac{1}{x}\left[\frac{4}{y^{4}}+\frac{16}{3 y^{5}}+\ldots\right] \\
& +\frac{1}{x^{2}}\left[-\frac{12}{y^{4}}-\frac{40}{3 y^{5}}+\ldots\right] \\
& +\frac{1}{x^{3}}\left[\frac{12}{y^{4}}+\frac{8}{y^{5}}+\ldots\right] \\
& +\frac{1}{x^{4}}\left[\frac{-4}{y^{4}}+\ldots\right] \\
& +\mathcal{O}\left(\frac{1}{x^{5} y^{4}}\right) \tag{8}
\end{align*}
$$

Note that, with $z=\tilde{z} e^{\mathbf{i} \theta}$, this leads to

$$
\begin{equation*}
S^{(2)}(\tilde{z}) \frac{(z-\tilde{z})^{2}}{2} \sim-\frac{1}{2} n^{\alpha} \theta^{2} \tag{9}
\end{equation*}
$$



Fig. 5: $G_{n}(z n) / G_{n}(\tilde{z}), n=500$, as function of $j$


Fig. 6: The quotient of the expression $\sqrt[8]{8}$ and $S^{(2)}(\tilde{z})$ as function of $j, n=500$

$$
\left.\begin{array}{rl}
S^{(3)}(\tilde{z}) & =\frac{1}{x^{2}}\left[-\frac{32}{y^{6}}+\ldots\right] \\
& +\frac{1}{x^{3}}\left[\frac{128}{y^{6}}+\ldots\right] \\
& +\frac{1}{x^{4}}\left[-\frac{192}{y^{6}}+\ldots\right] \\
& +\frac{1}{x^{5}}\left[\frac{128}{y^{6}}+\ldots\right] \\
& +\mathcal{O}\left(\frac{1}{x^{6} y^{6}}\right), \\
S^{(4)}(\tilde{z}) & =\frac{1}{x^{3}}\left[\frac{288}{y^{8}}+\ldots\right] \\
& +\frac{1}{x^{4}}\left[-\frac{1440}{y^{8}}+\ldots\right] \\
& +\frac{1}{x^{5}}\left[\frac{2880}{y^{8}}+\ldots\right] \\
& +\frac{1}{x^{6}}\left[-\frac{2880}{y^{8}}+\ldots\right] \\
& +\mathcal{O}\left(\frac{1}{x^{7} y^{8}}\right), \\
S^{(5)}(\tilde{z}) & =\frac{1}{x^{4}}\left[-\frac{3072}{y^{10}}+\ldots\right] \\
& +\frac{1}{x^{5}}\left[\frac{18432}{y^{10}}+\ldots\right] \\
& +\mathcal{O}\left(\frac{1}{x^{6} y^{10}}\right), \\
S^{(6)}(\tilde{z}) & =\frac{1}{x^{5}}\left[\frac{38400}{y^{12}}+\ldots\right] \\
& +\frac{1}{x^{6}}\left[\frac{268800}{y^{12}}+\ldots\right] \\
x^{7} y^{12}
\end{array}\right] .
$$

We proceed now as in Section 2 Again, the justification of the integration procedure is given in the Appendix. This leads to

$$
\tau=\frac{y^{2} \sqrt{x}}{2}\left[u a_{1}+\frac{u^{2} a_{2}}{x^{1 / 2}}+\frac{u^{3} a_{3}}{x}+\frac{u^{4} a_{4}}{x^{3 / 2}}+\frac{u^{5} a_{5}}{x^{2}}+\mathcal{O}\left(\frac{u^{6}}{x^{5 / 2}}\right)\right]
$$



Fig. 7: The quotient of the expression $\sqrt[8]{ }$ and $S^{(2)}(\tilde{z})$, as function of $j, n=500$. Restricted range, $\alpha \leq .84$ We give only $a_{1}$ :

$$
a_{1}=1-\frac{2}{3 y}-\frac{2}{9 y^{2}}+\ldots+\frac{1}{x}\left(\frac{3}{2}-\frac{4}{3 y}+\ldots\right)+\frac{1}{x^{2}}\left(\frac{15}{8}-\frac{7}{4 y}+\ldots\right)+\mathcal{O}\left(\frac{1}{x^{3}}\right)
$$

This leads to

$$
\frac{1}{\sqrt{2 \pi}} \int_{-\infty}^{\infty} e^{-u^{2} / 2} \tau^{\prime}(u) d u=\frac{y^{2} \sqrt{x}}{2} T_{4}
$$

with
$T_{4}=1-\frac{2}{3 y}-\frac{2}{9 y^{2}}+\ldots+\frac{1}{x}\left(\frac{5}{12}-\frac{11}{18 y}+\ldots\right)+\frac{1}{x^{2}}\left(\frac{73}{288}-\frac{133}{432 y}+\ldots\right)+\frac{1}{x^{3}}\left(\frac{721}{576}+\ldots\right)+\mathcal{O}\left(\frac{1}{x^{4}}\right)$.
Set

$$
T_{5}:=\frac{1}{\sqrt{2 \pi}} \frac{y^{2} \sqrt{x}}{2} e^{T_{1}} .
$$

This leads to

$$
\begin{equation*}
\left[z^{j}\right] \phi_{n}(z) \sim T_{5} T_{3} T_{4} \tag{10}
\end{equation*}
$$

We can of course combine $T_{3}$ and $T_{4}$ :

$$
\begin{aligned}
T_{6} & :=T_{3} T_{4}=1-\frac{3}{3 y}-\frac{1}{18 y^{2}}-\frac{1}{30 y^{3}}+\frac{17207}{3240 y^{4}}+\ldots+\frac{1}{x}\left(-\frac{1}{12}+\frac{1}{36 y}-\frac{35}{216 y^{2}}+\frac{15029}{3240 y^{3}}+\ldots\right) \\
& +\frac{1}{x^{2}}\left(\frac{1}{288}-\frac{1}{864 y}+\frac{3527}{5184 y^{2}}+\ldots\right)+\mathcal{O}\left(\frac{1}{x^{3}}\right)
\end{aligned}
$$



Fig. 8: The quotient $\left[z^{j}\right] \phi_{n}(z) / T_{8}$, two terms in $T_{4}$, as function of $j, n=500$


Fig. 9: The quotient $\left[z^{j}\right] \phi_{n}(z) / T_{8}$, three terms in $T_{4}$, as function of $j, n=500$

Let us consider the precision of our asymptotics.
The quality of asymptotic (6) is given in Figure 3 and 4 for $n=500$, and $x \in\left[\sqrt{n}, n^{0.9}\right]$ (first range) so that $y \in\left[n^{0.1}, \sqrt{n}\right]$. For some values of $j=n-x$, we show $\tilde{z} / z n$, where, as mentioned, $z n$ is the numerical solution of $S^{\prime}(z n)=0$. In the full range $j \in\left[n-n^{0.9}, n-\sqrt{n}\right]$, the precision is of order $10^{-5}$, in a restricted range, the precision is of order $10^{-6}$.

Also a comparison of $G_{n}(\tilde{z})$ and $G_{n}(z n)$ is given in Figure 5 , showing again a precision of order $10^{-6}$.
To check the quality of asymptotic (6), we give in Figure 6 the comparison between the expression (8) and $S^{(2)}(\tilde{z})$. The precision is of of order $10^{-2}$.
In a restricted range, given in Figure 7, the precision is of order $10^{-4} . \alpha \leq 0.84$ in this range.
We have made several experiments with (10), with $n$ up to 500 . The result is unsatisfactory, only values of $x$ of order $\sqrt{n}$ give reasonable results. Also using $e^{T_{2}}$ instead of $T_{3}$ does not improve the precision. Actually, only very large values of $n$ lead to good precision. So we turn to another formulation: instead of using $e^{T_{1}} T_{3}$ for $e^{S(\tilde{z})}$, we plug directly $\tilde{z}$ into $G_{n}(z)$, ie we set

$$
T_{7}=G_{n}(\tilde{z})
$$

leading to

$$
\left[z^{j}\right] \phi_{n}(z) \sim \frac{1}{\sqrt{2 \pi}} \frac{y^{2} \sqrt{x}}{2} T_{7} T_{4}=: T_{8} \text { say. }
$$

For $n=500$, using two and three terms in $T 4$, we give in Figures 8 and 9 the quotient $\left[z^{j}\right] \phi_{n}(z) / T_{8}$. The precision is of order $10^{-5}$.

## 4 Conclusion

Using an almost mechanized program in Maple, we have obtained some asymptotic expressions for Stirling numbers in central and non-central regions. We intend to use these techniques in other non-central ranges.

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## A Appendix. Justification of the integration procedure

## A. 1 The central region

We proceed as in Flajolet and Sedgewick [3], ch.VIII. We can choose here $\tilde{z}=1$. This leads, with $z=e^{\mathrm{i} \theta}$, to

$$
S(z) \sim S_{0}(z)+\mathcal{O}(\sqrt{\ln (n)} \theta)+\text { constant term }
$$

with

$$
\begin{aligned}
S_{0}(z) & =\sum_{k=0}^{n-1} \ln \left[e^{\mathbf{i} \theta}+k\right]-H_{n} \mathbf{i} \theta \\
& \sim \sum_{k=0}^{n-1} \frac{1}{1+k}\left[e^{\mathbf{i} \theta}-1\right]-\frac{1}{2} \sum_{k=0}^{n-1}\left[\frac{1}{1+k}\left[e^{\mathbf{i} \theta}-1\right]^{2}-H_{n} \mathbf{i} \theta+\mathcal{O}\left(\theta^{3}\right)\right. \\
& \sim H_{n}\left[e^{\mathbf{i} \theta}-1-\mathbf{i} \theta\right]+\mathcal{O}\left(\theta^{2}\right)
\end{aligned}
$$

Set

$$
h(\theta):=e^{\mathbf{i} \theta}-1-\mathbf{i} \theta
$$

We have

$$
h(\theta) \sim-\frac{\theta^{2}}{2}
$$

which conforms to (3).
The function $h(\theta)$ is the same as in [3], Ex.VIII.3, which proves the validity of our integration procedure: we use here $H_{n} \sim \ln (n)$ instead of $n$. The complete asymptotic expansion is justified as in [3], Ex.VIII.4.

## A. 2 The non-central region

We choose here $\tilde{z}=\frac{n y}{2}=\frac{n^{2-\alpha}}{2}:=\delta$, say. We have

$$
\begin{aligned}
& \frac{1}{2}<\alpha<1, \\
& n^{\alpha}=\frac{n^{2}}{2 \delta}, \\
& n^{2} \gg \delta \gg n \gg n^{\alpha} .
\end{aligned}
$$

Set $z=\delta e^{\mathbf{i} \theta}$, this leads, with Euler-Maclaurin formula, with the first correction (the other corrections are negligible), to

$$
\begin{aligned}
S(z) & \sim \sum_{k=0}^{n-1} \ln \left[\delta e^{\mathbf{i} \theta}+k\right]-\left(n-n^{\alpha}\right) \mathbf{i} \theta-\left(n-n^{\alpha}\right) \ln (\delta) \\
& \sim \ln \left[n+\delta e^{\mathbf{i} \theta}\right]\left[n+\delta e^{\mathbf{i} \theta}\right]-n-\delta e^{\mathbf{i} \theta} \ln \left[\delta e^{\mathbf{i} \theta}\right]-\left[n-\frac{n^{2}}{2 \delta}\right](\mathbf{i} \theta+\ln (\delta))-\frac{1}{2} \ln \left[n+\delta e^{\mathbf{i} \theta}\right]+\frac{1}{2} \ln \left[\delta e^{\mathbf{i} \theta}\right]
\end{aligned}
$$

Set now $n=\rho \delta, \rho=2 n^{\alpha-1} \ll 1$ and expand wrt $\rho$. This gives

$$
\begin{aligned}
S(z) & \sim \rho\left[-\frac{1}{2} e^{-\mathbf{i} \theta}\right] \\
& +\rho^{2}\left[\delta \frac{1+\mathbf{i} \theta e^{\mathbf{i} \theta}}{2 e^{\mathbf{i} \theta}}+\frac{1}{4} e^{-2 \mathbf{i} \theta}+\frac{1}{2} \delta \ln (\delta)\right] \\
& +\rho^{3}\left[-\frac{\delta}{6} e^{-2 \mathbf{i} \theta}-\frac{1}{6} e^{-3 \mathbf{i} \theta}\right] \\
& +\mathcal{O}\left(\delta \rho^{4}\right) .
\end{aligned}
$$

Note that the dominant constant contribution is given by $\frac{1}{2} \rho^{2} \delta \ln (\delta)=(2-\alpha) n^{\alpha} \ln (n)$, which conforms to 7 . The first term gives a variable part $\mathcal{O}(\rho)$. The second term gives a variable part $2 n^{\alpha} h(\theta)+\mathcal{O}\left(\rho^{2}\right)$, with

$$
h(\theta):=\frac{1+\mathbf{i} \theta e^{\mathbf{i} \theta}}{2 e^{\mathbf{i} \theta}} .
$$

The third term give $\mathcal{O}\left(n^{2 \alpha-1}\right) \ll n^{\alpha}$. Note that $2 n^{\alpha} h(\theta) \sim-\frac{1}{2} n^{\alpha} \theta^{2}$, which conforms to 99. The function $\left|e^{h(\theta)}\right|=e^{\cos (\theta) / 2}$ is unimodal with peak at 0 and $h(0)=1 / 2$. Let us introduce a splitting value $\theta_{0}$ such that $n^{\alpha} \theta_{0}^{2} \rightarrow \infty, n^{\alpha} \theta_{0}^{3} \rightarrow 0, n \rightarrow \infty$. For instance, we choose $\theta_{0}=n^{\beta}, \beta=-\frac{5 \alpha}{12}$. By unimodality property of the cosine, the tail integral

$$
K_{n}^{(1)}:=\int_{\theta_{0}}^{2 \pi-\theta_{0}} e^{2 n^{\alpha}(h(\theta)-1 / 2)} d \theta
$$

is such that

$$
\left|K_{n}^{(1)}\right|=\mathcal{O}\left(e^{n^{\alpha}\left[\cos \left(\theta_{0}\right)-1\right]}\right)=\mathcal{O}\left(e^{-C n^{\alpha / 6}}\right)
$$

for some $C>0$. The tail integral is exponentially small.
As $h(\theta) \sim-\frac{\theta^{2}}{4}$, the central approximation and the tail completion are immediate.

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